

Stars: beyond the Main Sequence

Main-sequence

- 1) first and dominant stage in life,
- 2) $H \rightarrow He$ in the core (subtle rise of L)
- 3) >80% stars in MS

Low-mass stars

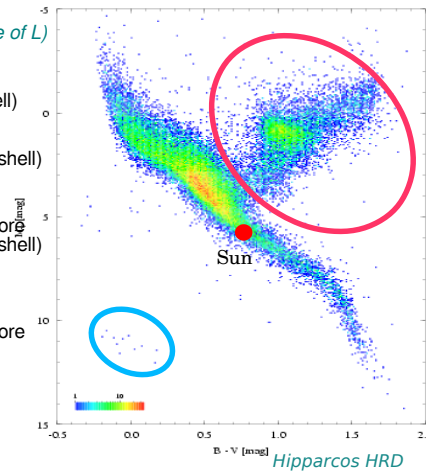
- 1) Ascend **Red Giant Branch** (H shell)
- 2) **Red clump/Horizontal branch** (He \rightarrow C in core, H shell)
- 3) **Asymptotic Giant Branch** (H, He shell)
- 4) Env. ejection \rightarrow **White Dwarf**

Intermediate-mass stars

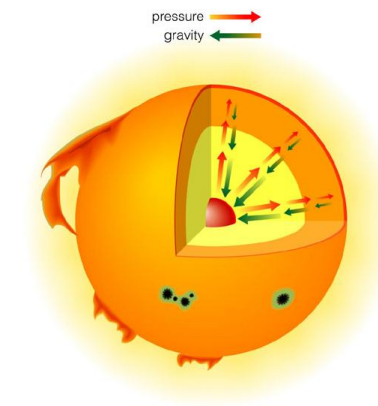
- 1) Expand at $L \sim \text{const}$, ignite He in core
- 2) **Asymptotic Giant Branch** (H, He shell)
Some: "super-AGB" (H, He, C shell)
- 3) Env. ejection \rightarrow **White Dwarf**

High-mass stars

- 1) Expand at $L \sim \text{const}$, ignite He in core
- 2) Fusion in stages, up to Si \rightarrow Fe
- 3) Core collapse \rightarrow **Neutron star**,
Black hole



Star's life: Protracted battle with gravity



- ALWAYS** \Rightarrow To support weight:
 \Rightarrow need high pressure
- MOSTLY** \Rightarrow need high temperature
 \Rightarrow will loose energy
 \Rightarrow need energy source:
- Gravitational contraction
- Nuclear fusion

Ultimately,
Can something else than thermal pressure balance gravity?

Degeneracy for a White Dwarf or degenerate core

using wave-particle duality

$$M = 1 M_{\odot}, R = 6000 \text{ km, electron mean spacing } d \sim \frac{1}{n_e^{1/3}} \sim 10^{-12} \text{ m}$$

$$\text{De Broglie wavelength } \lambda \sim \frac{h}{p} \sim \frac{h}{m_e v} \sim 10^{-12} \text{ m} \left(\frac{10^9 \text{ K}}{T} \right)^{1/2}$$

\rightarrow typically $\lambda > d$ in white dwarfs

1. Wave nature of electrons becomes relevant
2. Electrons are fermion: they do not like to be close; this exclusion provides pressure

electron degeneracy pressure: $P \propto n_e^{5/3} \propto \rho^{5/3} \rightarrow R \propto M^{-1/3}$
Sirius B ($1 M_{\odot}$) smaller than typical white dwarfs ($0.6 M_{\odot}$)

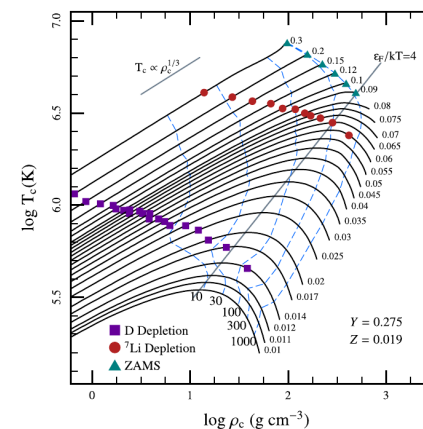
(compare with MS: $P = nkT$ plus $T \sim \text{constant} \rightarrow R \propto M$)

Can $M \rightarrow \infty$? No, would need $v > c \rightarrow$ Chandrasekhar limit: $M < 1.4 M_{\odot}$

White dwarfs are supported against gravity by pressure from degenerate electrons (ions irrelevant). This pressure depends only on density, not on temperature.
-- also centre of evolved low-mass stars, brown-dwarfs, Jupiter, metals...

Brown dwarf ($M < 0.08 M_{\odot}$)

no H fusion \rightarrow never becomes star



Ultimately,
Can something else than thermal pressure balance gravity?

Below $\sim 0.08 M_{\odot}$,
objects become degenerate
before becoming hot enough
to ignite H fusion
 \rightarrow **Brown dwarfs**

Below $\sim 0.013 M_{\odot}$,
no fusion at all (not even D)
 \rightarrow **Planet?**

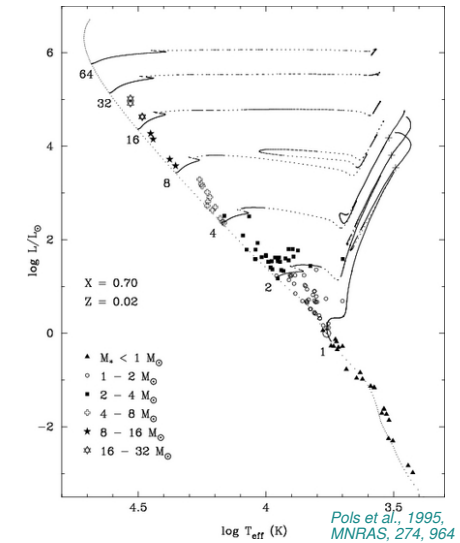
Low-mass star ($M < 8 M_{\odot}$)

gradual exhaustion of H & He
ends due to mass loss

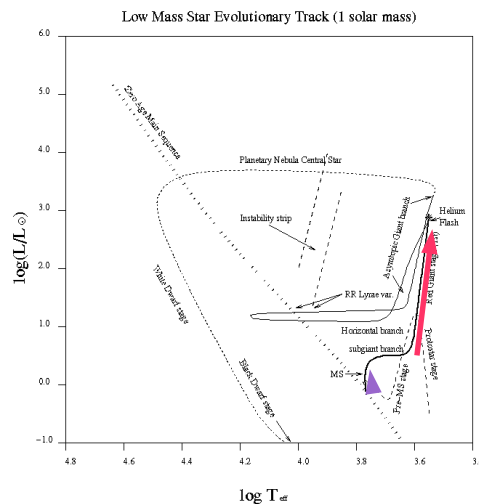
- Main sequence** Core H burning; for $\sim 1 M_{\odot}$, $dM/dt \sim 10^{-14} M_{\odot}/\text{yr}$, live $\sim 10^{10}$ yr $(M/M_{\odot})^{-3}$
- Red Giant** Dense core + tenuous envelope; $\sim 100 R_{\odot}$, $\sim 10^8$ yr, shell H burning, mass loss $dM/dt \sim 10^{-7} M_{\odot}/\text{yr}$; envelope convective, core shrinks (*why?*)
- Red clump/Horizontal Branch** $T_{\text{core}} \sim 10^8$ K, He fuses in core, but H shell dominates L; $\sim 10\%$ of main-sequence lifetime; move to blue if metal-poor
- Asymptotic Giant** C/O core + extremely tenuous envelope, $\sim 600 R_{\odot}$; $\sim 10^6$ yr, $10^5 L_{\odot}$, shell He/H burning, He shell flashes; end of AGB: $dM/dt \sim 10^{-5} M_{\odot}/\text{yr}$, removes whole envelope
- White dwarf** Dense C/O, $\sim 0.6 M_{\odot}$, $\sim 0.01 R_{\odot}$, e^- degeneracy, cools from $\sim 10^5$ K to few 10^3 K in $\sim 10^{10}$ yr
- +planetary nebula Few $0.1 M_{\odot}$, expansion speed ~ 20 km/s, fluoresces for $\sim 10^4$ yr

Net effect of a star's life: mass loss \rightarrow metal enrichment of the interstellar medium

Evolutionary tracks:
what happens depends on mass



Pols et al., 1995, MNRAS, 274, 964



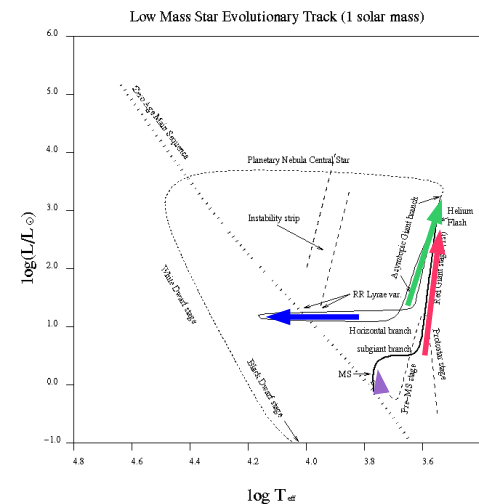
$$L = 4 \pi R^2 \sigma T^4$$

Main-sequence
Earlier: lower L; paradox of the 'faint young sun'

Red giant branch
Shell H burning around degenerate He core. Dense core sets properties of shell; e.g., $T \sim M_{\text{core}}/R_{\text{core}}$. Thus, core determines L. Star *has to* transport this \rightarrow large radius increase.

As core increases, so does T (and L drastically; why?)
Until He ignites \rightarrow He flash

From <http://www.ngcsu.edu/Academic/Sciences/physics/jones/ast1020home/evoltracks.htm>



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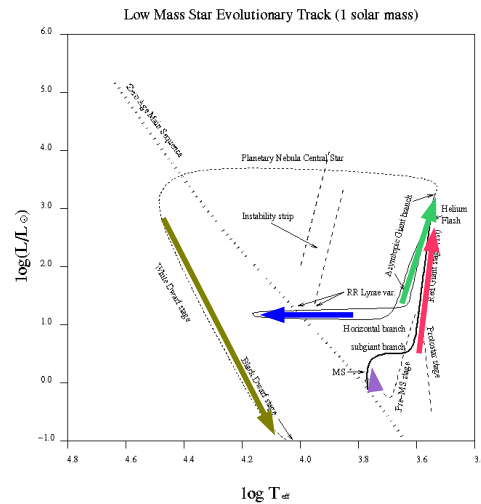
Red giant branch
Shell H burning
Big, luminous (cook planets)

Red clump/Horizontal branch
Core He burning
Core bigger $\rightarrow T_{\text{shell}}$ lower
 \rightarrow less luminous, less big
Hotter if metal-poor

Asymptotic Giant branch
Shell H & He burning
Luminosity increases with increasing core mass, star becomes very big.

Ends when L so high that envelope blown off

From <http://www.ngcsu.edu/Academic/Sciences/physics/jones/ast1020home/evoltracks.htm>



$$L = 4 \pi R^2 \sigma T^4$$

Main-sequence

Earlier: lower L; paradox of the 'faint young sun'

Red giant branch

Shell H burning
Big, luminous (cook planets)

Red clump/Horizontal branch

Core He burning
Bit less big & luminous
Hotter if metal-poor

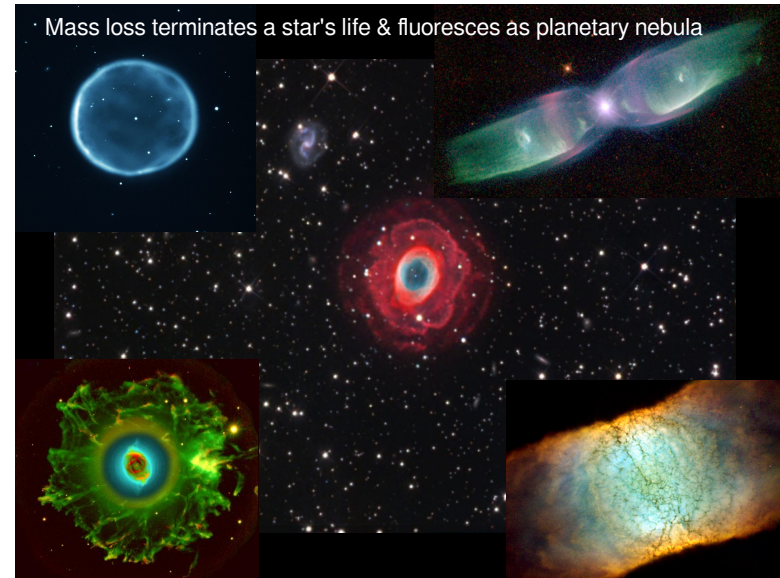
Asymptotic Giant branch

Shell H & He burning
Very big (engulfing planets)

White Dwarf branch

Naked stellar core
Born very hot
Cools off slowly

From <http://www.ngcsu.edu/Academic/Sciences/physics/jones/ast1020home/evoltracks.htm>



Mass loss terminates a star's life & fluoresces as planetary nebula

Extra Notes: Planetary Nebula

1) what are they?

Transient fluorescence around new-born, hot white dwarfs

Live $\sim 10^4$ yr

Where is the gas from?

Slow expansion (20 km/s)

2) what are the colors?

'False colors,' montage of narrow-band images

e.g., red \rightarrow N I

green \rightarrow H α

blue \rightarrow O III

3) why aren't they all spherical?

Rotation?

Binaries?



Extra Notes: degeneracy pressure

1. Wave nature of electrons becomes relevant

2. Electrons are fermion: they do not like to be close; this exclusion provides pressure

Pauli's exclusion principle: $(\Delta x_i)^3 (\Delta p_i)^3 \sim \hbar^3$

Fermi momentum: $p_F \sim \frac{1}{\Delta x_i} \sim n_e^{1/3}$

Fermi energy: $E_F = \frac{1}{2} m_e v_F^2 = \frac{1}{2} \frac{p_F^2}{m_e} \propto n_e^{2/3}$

Pressure: $P \propto n_e E_F \propto n_e^{5/3} \propto \rho^{5/3}$

$\rightarrow R \propto M^{-1/3}$

(compare with MS: $P = nkT$ plus $T \sim \text{constant} \rightarrow R \propto M$)

Can $M \rightarrow \infty$? No, would need $v > c \rightarrow$ Chandrasekhar limit: $M < 1.4 M_\odot$