

$$H\ddot{E} \quad \frac{dP}{dr} = -\frac{GM_r}{r^2} \rho \Rightarrow \frac{P_c}{R} \approx \frac{GM}{R^2} \rho \approx \frac{GM^2}{R^5}$$

$$P_c \approx \frac{GM^2}{R^4}$$

disturbance $\Rightarrow R$ a little smaller \rightarrow homologously

to be stable: real $P \propto R^{-4}$ or steeper

what we know: $\rho \propto R^{-3}$

ideal gas $P \propto \rho T$ isothermal $P \propto \rho \propto R^{-3}$

not steep enough \Rightarrow unstable

adiabatic $P \propto \rho^\gamma \leftarrow \begin{matrix} \text{adiabatic} \\ \text{index} \end{matrix} \propto R^{-3\gamma}$

stability requires $\boxed{\gamma > \frac{4}{3}}$
 ideal gas $\gamma = \frac{5}{3} \Rightarrow$ stable

dominated by radiation pressure

$$P = \frac{1}{3} a T^4$$

compress it $n_\gamma \propto R^{-3}$
 mean energy of photons $\epsilon_\gamma \propto R^{-1}$ } $P \propto R^{-4}$
 marginally stable

$$P_{\text{true}} = P_{\text{rad}} + P_{\text{ideal}}$$

need to have some suitable average of γ
 (can have small isothermal core)

how to get $\gamma < \frac{4}{3}$?

Ionisation zone \rightarrow mix of H^0, p^+, e^-
 compress a little $\rightarrow T \uparrow \Rightarrow$ ionize more H^0
 \rightarrow costs energy $\Rightarrow T \downarrow$ a little $\Rightarrow T$ wind up as fast

γ lower $\uparrow \uparrow$

Time scales

Dynamical time scale

$$g = \frac{GM}{R^2}$$

not in HE: $\rho \frac{\partial v}{\partial t} = \rho \frac{\partial^2 r}{\partial t^2} = -\nabla P + \rho \nabla \phi = -\frac{\partial P}{\partial r} - \rho g$

ignore ρ : $\rho \frac{\partial^2 r}{\partial t^2} = -\rho g \Rightarrow \frac{R}{t_{ff}^2} \approx -g$

$$\Rightarrow t_{ff} \approx \sqrt{\frac{R}{g}} \approx \frac{1}{\sqrt{G\rho}} \quad \rho_{\odot} \approx 1 \text{ g/cm}^3 \Rightarrow t_{ff} \approx 1 \text{ hr}$$

ignore gravity $\rho \frac{\partial^2 r}{\partial t^2} = -\frac{\partial P}{\partial r}$

$$\rho \frac{R}{t_{expl}^2} \approx \frac{P}{R} \Rightarrow t_{expl} = R \sqrt{\frac{\rho}{P}} \approx R/c_s$$

$$\rho \approx \frac{GM^2}{R^4} \Rightarrow t_{expl} = \frac{1}{\sqrt{G\rho}}$$

HE: fast \Rightarrow safe to assume stars are in HE

Thermal equilibrium

$$t_{th} = \frac{E_{th}}{L} \approx \frac{E_{pot}}{L} \approx \frac{GM^2}{RL}$$

$N_{th} \approx \frac{M}{\mu m_H}$ at

no TE in a small spot, over some distance d

\odot : few 10^7 yr

$$F = -\frac{1}{3} v_{lmp} \nabla U \approx \frac{1}{3} \frac{c}{\kappa \rho} \frac{u}{d}$$

much colder



$$E_{adj} = \frac{Ud^3}{Fd^2} = \frac{E_{blob}}{L_{adj}/blob}$$

$$= \frac{3Ud^3}{v_{lmp} U d} \approx \frac{3d^2}{v_{lmp}} \approx \left(\frac{v_{lmp}}{c}\right) \left(\frac{d}{v_{lmp}}\right)^2$$

Nuclear time scale

$$\eta \approx \frac{1 \text{ MeV/nucleon}}{m_{nucleon} c^2} \approx 0.1\% \approx 0.001 \quad t_{nuc} = \frac{E_{nuc}}{L} = \frac{\eta M_{fuel} c^2}{L}$$

in random walk $t_{step} \downarrow N_{steps}$

Polytropes

$$P \propto \rho^\gamma$$

adiabatic expansion
contraction

$$\frac{dM_r}{dr} = 4\pi r^2 \rho$$

$$\frac{dP}{dr} = -\frac{GM_r}{r^2} \rho$$

$$P = k \rho^\gamma$$

polytropic equation
of state

normally $P = P(\rho, T, \text{abundances})$

$$P = k \rho^{1 + \frac{1}{n}} \leftarrow \text{polytropic index}$$

$$\gamma = \frac{5}{3} \Rightarrow n = 1.5$$

$$\gamma = \frac{4}{3} \Rightarrow n = 3$$

$$P_c = \text{const}_1 \frac{GM^2}{R^4}$$

$$\rho_c = \text{const}_2 \frac{M}{\frac{4}{3}\pi R^3}$$